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## THE STRUCTURAL PECULIARITIES OF A CHIRAL LIQUID CRYSTAL IN RAMAN SPECTRA CONTOURS

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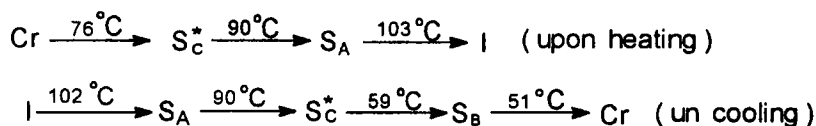
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**Abstract** Fourier transform ( FT ) Raman spectroscopy is shown to be a very powerful and valuable technique for the analysis of liquid crystal materials. By temperature-variable FT-Raman spectroscopy, we found a new crystalline phase of a chiral liquid crystal compound and studied the intermolecular interaction of the compound in its smectic phases (  $S_A, S_C^*, S_B$  ).

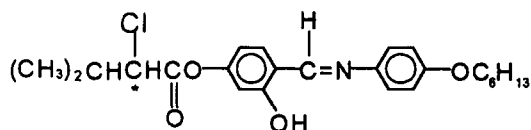
### INTRODUCTION

Generally speaking, the properties of liquid crystals depend on the structure of liquid crystal molecules, so the study of the relationship between molecular structure and the properties of liquid crystals is very important and interesting. Vibrational spectroscopy has proved to be useful for studying the structure and physical properties of liquid crystals. Temperature-Variable FT-Raman Spectroscopy is one of the most powerful tools for studying the molecular environment of liquid crystals.<sup>1,2</sup>

In our previous paper,<sup>3</sup> we presented the synthesis of this new chiral liquid crystal compound. And the differential scanning calorimetry (DSC), optical polarizing microscopy and X-ray diffraction measurements were used to study the phase transition temperatures and behaviour. The results reveal that this compound shows the following transition<sup>3</sup>:



In this paper, we reported the structural studies of this compound by temperature-variable FT-Raman spectroscopy. Its chemical structure is :



On cooling, the results obtained by Raman spectroscopy is corresponding well with the conclusions of DSC and optical polarizing microscopy<sup>3</sup>, the investigated compound exhibited  $S_A$ ,  $S_C^*$  and  $S_B$  phases. But upon heating, we observed two crystalline phases by Raman spectroscopy.

### EXPERIMENTAL

The sample was analysed in the solid state using a Bruker RFS-100 FT-Raman instrument. The sample was scanned 500 times in order to obtain a good signal-to-noise. The output laser power is 50mw. Sample was mounted as powder in a hot stage. The rate of temperature changing is 2°C/min in both heating and cooling processes. In this cycle before Raman spectra registration the sample was kept at each temperature point for 15 minutes.

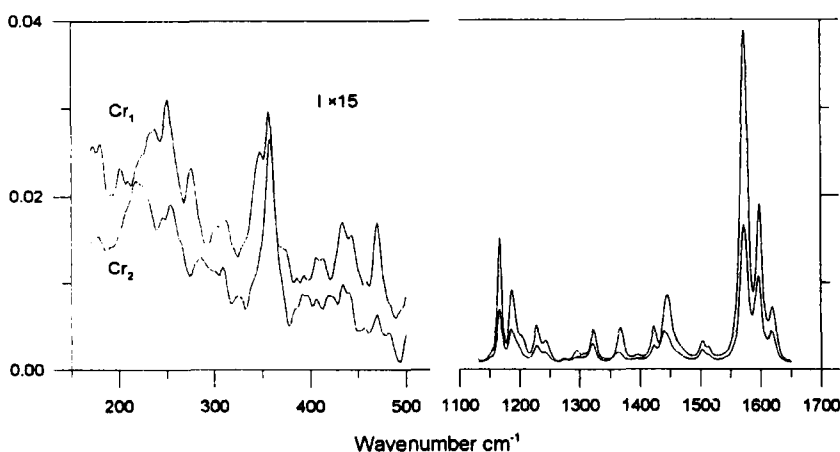


FIGURE 1 FT-Raman spectra of  $Cr_1$  phase (above curve) and  $Cr_2$  phase (below curve) in 150-500  $cm^{-1}$  and 1130-1650  $cm^{-1}$  regions.

## RESULTS AND DISCUSSION

When heating the new sample, below the Cr-Sc\* transition temperature we observed a clearly change in Raman spectra, see Figure 1. The most obvious changes are the two bands at  $1446\text{cm}^{-1}$  and  $1367\text{cm}^{-1}$ . The band  $1446\text{cm}^{-1}$  which is assigned to the asymmetric  $\text{CH}_3$  deformation and the  $\text{CH}_2$  scissors deformation shifted to  $1441\text{cm}^{-1}$ , the band  $1367\text{cm}^{-1}$  the symmetric  $\text{CH}_3$  deformation shifted to  $1362\text{cm}^{-1}$  and became broad in shape and decreasing in intensity.<sup>4</sup> These changes indicate that the molecule vibrations and the interactions between molecules are all changed. But in the studies of phase behaviour by DSC and optical polarizing microscopy<sup>3</sup>, we don't find this kind of phase transition. So we think it does not a crystalline-liquid crystal phase transition, this transition occurs between two crystalline phases, i.e.  $\text{Cr}_1 \rightarrow \text{Cr}_2$ . Continuing heating, one after another we observed the structure peculiarities of the  $\text{S}_\text{C}^*$ ,  $\text{S}_\text{A}$  and isotropic (I), respectively. So now the compound shows the following transition upon heating:  $\text{Cr}_1 \rightarrow \text{Cr}_2 \rightarrow \text{S}_\text{C}^* \rightarrow \text{S}_\text{A} \rightarrow \text{I}$ .

When cooling the compound from the isotropic phase, by Raman spectroscopy we obtained the structure peculiarities of the following phases transitions:  $\text{I} \rightarrow \text{S}_\text{A} \rightarrow \text{S}_\text{C}^* \rightarrow \text{S}_\text{B} \rightarrow \text{Cr}_2$ . Generally speaking, in smectic A phase the lateral distribution of the molecules within each layer is however random, and the molecules in the  $\text{S}_\text{A}$  layers are able to rotate freely about their long axes.<sup>5</sup> In Sc phase there is no long range correlation between the layers, the constituent molecules are randomly packed within the layers but the long axes of the molecules are tilted with respect to the layer normal<sup>5</sup> (For this compound the tilt angle in the  $\text{S}_\text{C}^*$  phase is in the range of  $2^\circ$  to  $30^\circ$ ).<sup>3</sup> The molecules rotate about their long axes. In  $\text{S}_\text{B}$  phase the molecules are still arranged in layers, but the molecular long axes are shown in a hexagonally close-packed array. So the rotation of a individual molecule is confined, thus the molecular rotation must be completely cooperatively.<sup>5</sup> That is to say, on cooling from  $\text{S}_\text{A}$  to  $\text{S}_\text{C}^*$  then to  $\text{S}_\text{B}$  and  $\text{Cr}_2$ , the intermolecular interaction of molecules is becoming stronger and stronger and the freedom degree of molecular motions is becoming smaller and smaller.

Now let us look at the Raman spectra of the investigated compound in the  $200\text{--}450\text{cm}^{-1}$  region, see Figure 2. The bands in this region are associated with a group

motions of the hexyl side groups, most probably the all trans accordian band.<sup>6</sup> These bands originated from the CCC deformations containing some C-C stretching, it may lead to chain stretch out and draw back.<sup>4</sup> Following the decreasing of the freedom

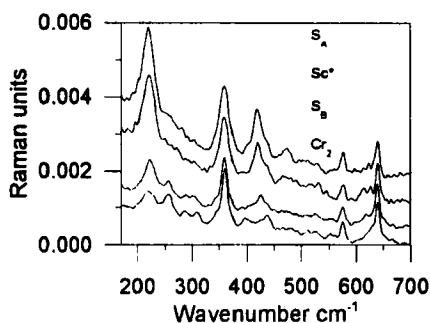


FIGURE 2 FT-Raman spectra of  $S_A$ ,  $S_C^*$ ,  $S_B$  and  $C_2$  phases in region  $170\text{--}700\text{cm}^{-1}$ .

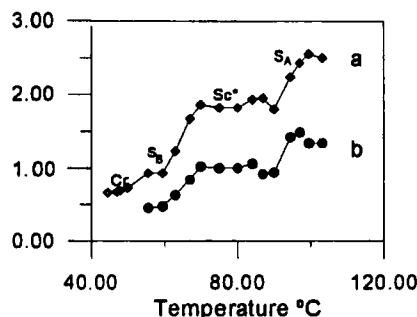


FIGURE 3 Temperature dependencies of relative integral value:  $a = I_{218}/I_{639}$ ,  $b = I_{418}/I_{639}$ .

degree of the molecular motions, we observed the changes in the intensities of these bands. Figure 3 illustrated the temperature dependencies of integral values of bands  $218\text{cm}^{-1}$  and  $418\text{cm}^{-1}$ . In the whole thermal circulation the band at  $639\text{cm}^{-1}$  which we assigned to C-Cl stretching model<sup>6</sup> remains no change. We use this band as a internal sign, so the integral values in Figure 3 are not the absolute values of those bands, there are only the relative values, i.e.  $I_{218}/I_{639}$  and  $I_{418}/I_{639}$ . From Figure 3 we can say that the bands at  $218\text{cm}^{-1}$  and  $418\text{cm}^{-1}$  are the signs for the freedom degree of the molecular motions. The shifts and the appearances of new bands at the liquid crystal-crystal transition indicate that the environment of the molecule is changed, i.e. the intermolecular interaction is modified.

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